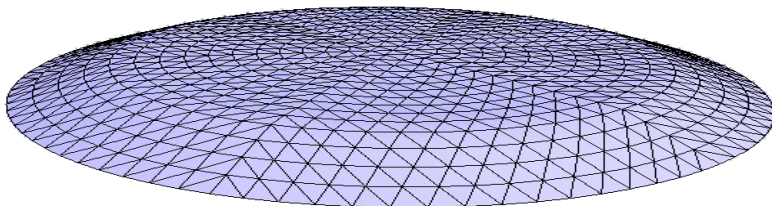




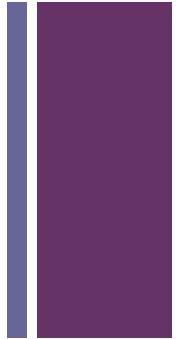
Laplace-Beltrami Eigenstuff Part 1 - Background

Martin Reuter – reuter@mit.edu

Mass. General Hospital, Harvard Medical, MIT

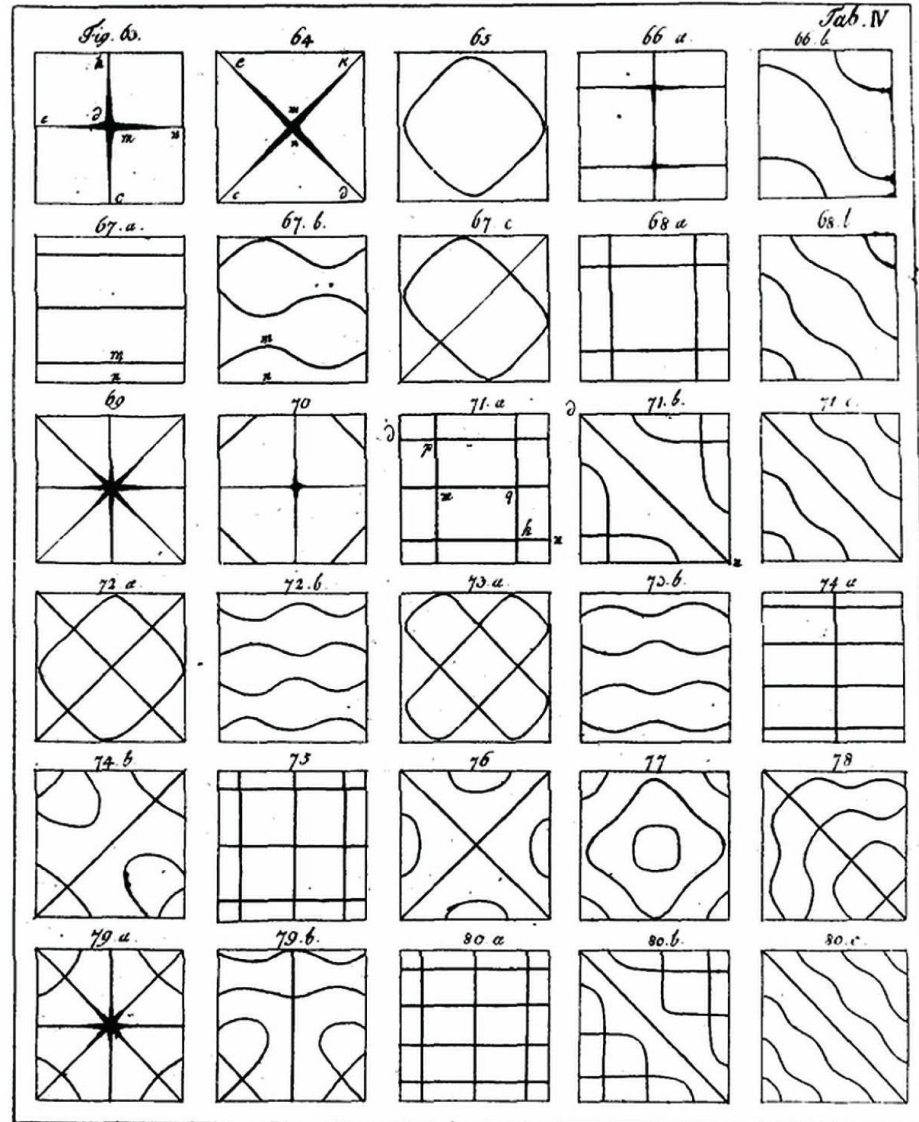


+ Sound and Shape



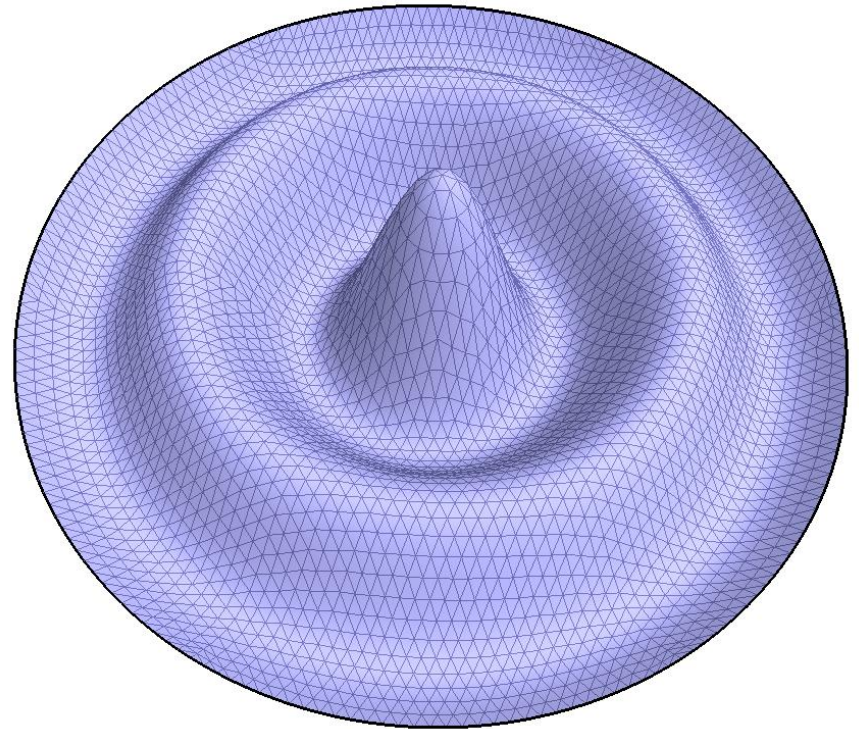
+ Chladni's strange patterns

- vibration of plates
- used a violin bow
- discovered sound patterns by spreading sand on the plates
- 1809 invited by Napoleon
- who held out price for mathematical explanation
- “Entdeckungen über die Theorie des Klanges” (Discoveries concerning the theory of music), Chladni, 1787



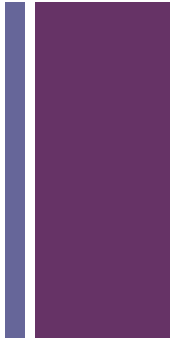
+ Can one “hear” Shape?

- First asked by Bers, then paper by Kac 1966, idea dates back to Weyl 1911 (at least).
- The sound (eigen-frequencies) of a drum depend on its shape.
- This spectrum can be numerically computed if the shape is known.
- E.g., no other shape has the same spectrum as a disk.
- Can the shape be computed from the spectrum?





Contributions



Shape Analysis: Reuter, Wolter, Peinecke [SPM05],
[JCAD06](most cited paper award 09) and Patent Appl.

- Introduced Laplace-Beltrami Op. for Non-Rigid Shape Analysis.
- Cubic FEM to obtain accurate solutions for surfaces and solids.
- Before a mesh Laplace (simplified linear FEM) has been used for parametrization, smoothing, mesh compression

Image Recognition: Peinecke et al [JCAD07] (Mass Density LBO)

Neuroscience Applications:

- Statistical morphometric studies of brain structures (eigenvalues), Niethammer, Reuter, Shenton, Bioux.. [MICCAI07], Reuter.. [CW08]
- Topological studies of eigenfunctions, Reuter.. [CAD09] (invited)

Segmentation: Reuter..(IMATI, Genova) [SMI09] [IJCV09]

Correspondence: Reuter [IJCV09]

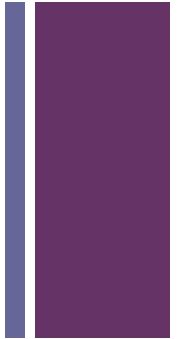
+ Why is the Laplace Operator so interesting?

involved in many PDE's:

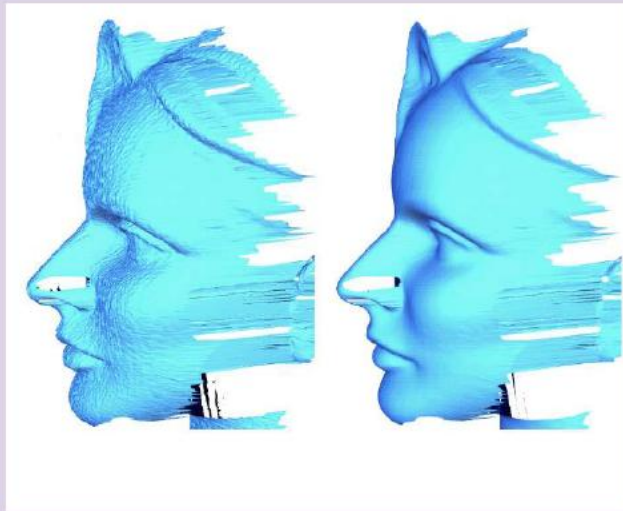
- Poisson equation: $\Delta f = h$
- Laplace equation: $\Delta f = 0$
- Heat equation: $\frac{\partial f}{\partial t} - k\Delta f = 0$
- Wave equation: $\frac{\partial^2 f}{\partial t^2} - k\Delta f = 0$
- Helmholtz equation: $\Delta f = -\lambda f$

=> Many properties carry over to potential applications from Physics ...

+ Applications



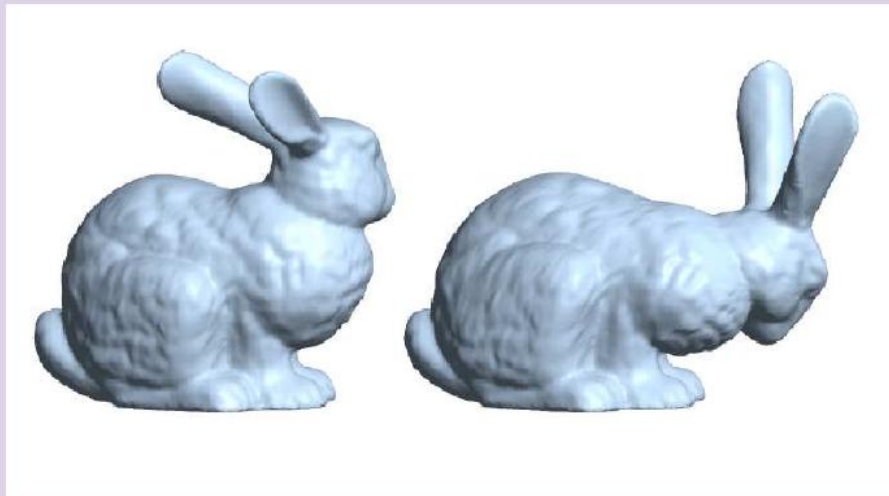
Mesh Smoothing



[Desbrun et al 1999, etc]

+ Applications

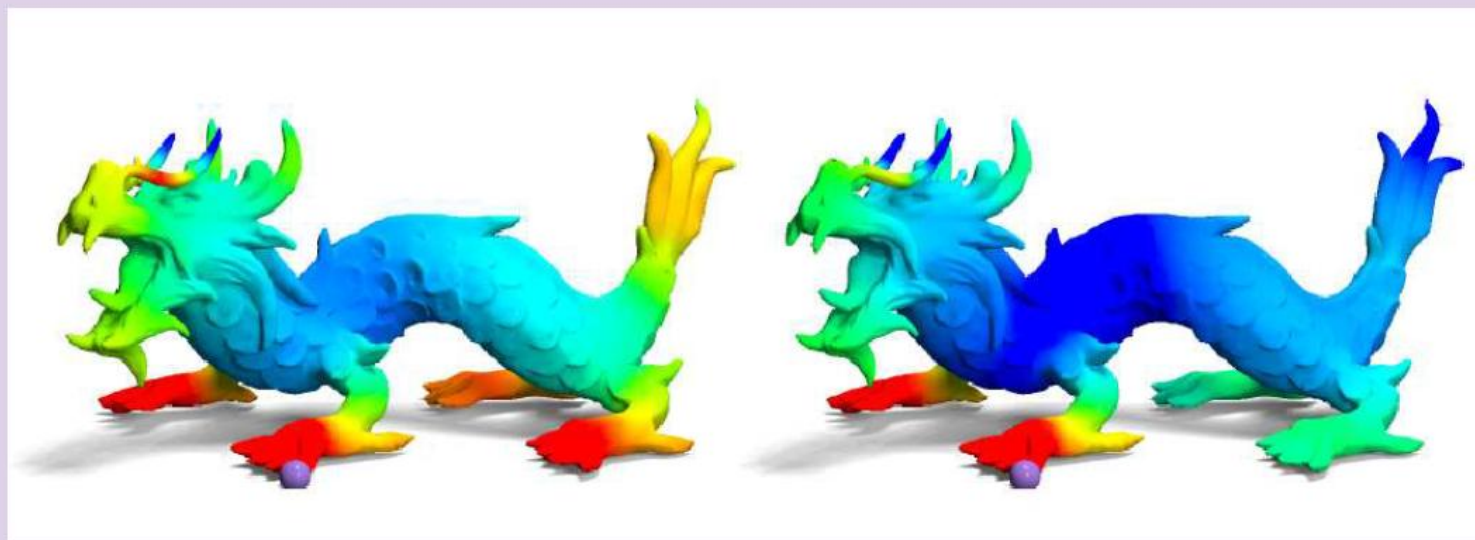
Mesh Editing



[Zhou et al 2005, Lipman et al 2005, etc]

+ Applications

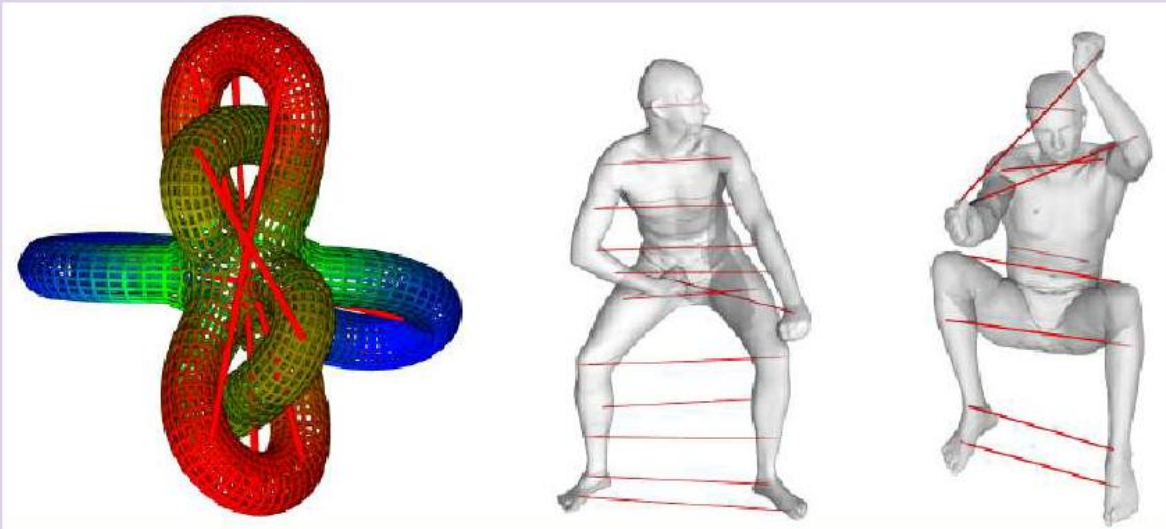
Heat Kernel Signature



[Sun, Ovsjanikov, and Guibas 2008, etc]

+ Applications

Shape Analysis



[Ovsjanikov, Sun and Guibas 2008, etc]

+ Laplace Spectrum

Definition

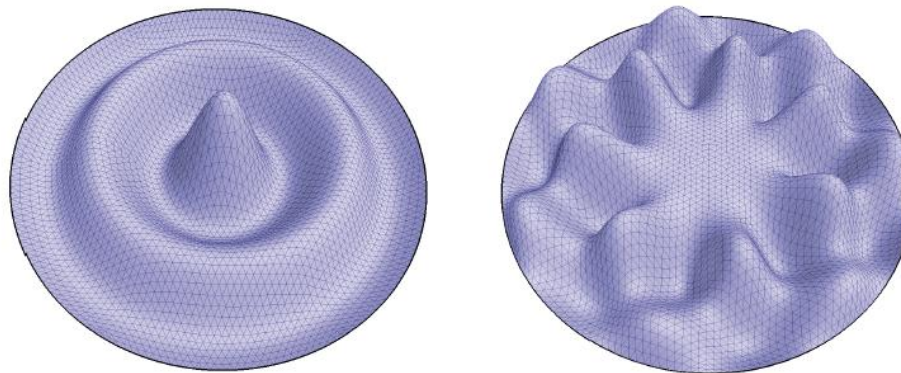
Helmholtz Equation (Laplacian Eigenvalue Problem):

$$\Delta f = -\lambda f, \quad f : M \rightarrow \mathbb{R}$$

Solution: Eigenfunctions f_i with corresponding family of eigenvalues (**Spectrum**):

$$0 \leq \lambda_1 \leq \lambda_2 \leq \cdots \uparrow +\infty$$

Here Laplace-Beltrami Operator: $\Delta f := \operatorname{div}(\operatorname{grad} f)$



+ Boundary Conditions

Laplace-Beltrami Spectrum for Manifolds with Boundary:

Dirichlet Boundary Condition

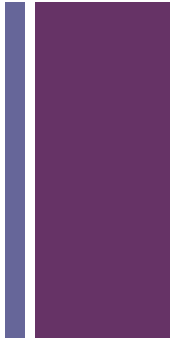
Function is fixed $f \equiv 0$ on the boundary of M

Neumann Boundary Condition

Derivative in normal direction is fixed $\frac{\partial f}{\partial n} \equiv 0$ on the boundary of M



1-Dimensional Helmholtz



$$\Delta f = -\lambda f \quad \Omega = \{x : 0 \leq x \leq a\}$$

$$\Leftrightarrow f'' = -\lambda f \Leftrightarrow f'' + \lambda f = 0$$

Neumann Boundary Condition :

$$f'(0) = 0 \quad \text{and} \quad f'(a) = 0$$

$$\Rightarrow f(x) = c \cdot \cos\left(\frac{n\pi}{a}x\right) \quad n = 0, 1, 2, \dots$$

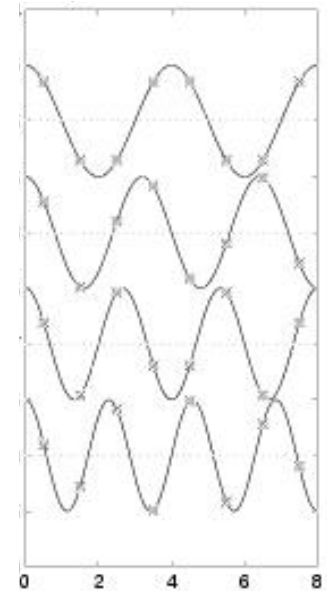
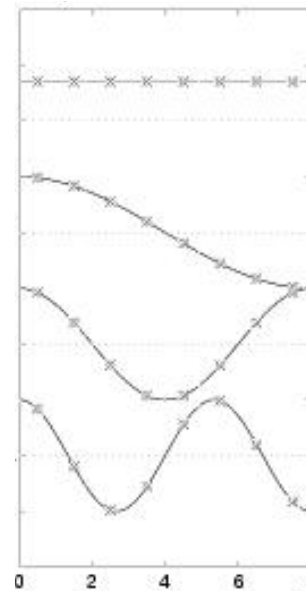
$$\Rightarrow \lambda = \left(\frac{n\pi}{a}\right)^2$$

Dirichlet Boundary Condition :

$$f(0) = 0 \quad \text{and} \quad f(a) = 0$$

$$\Rightarrow f(x) = c \sin(n\pi x / a) \quad n = 1, 2, 3, \dots$$

$$\Rightarrow \lambda = (n\pi / a)^2$$



+ 2D Rectangle



- Side lengths a and b (Neumann Boundary Condition)

- Separation of variables leads to Eigenfunctions:

$$f(x, y) = c \cos\left(\frac{m\pi}{a}x\right) \cos\left(\frac{n\pi}{b}y\right) \quad m, n = 0, 1, 2, 3, \dots$$

- Eigenvalues:

$$\lambda = \pi^2 \left(\frac{m^2}{a^2} + \frac{n^2}{b^2} \right)$$

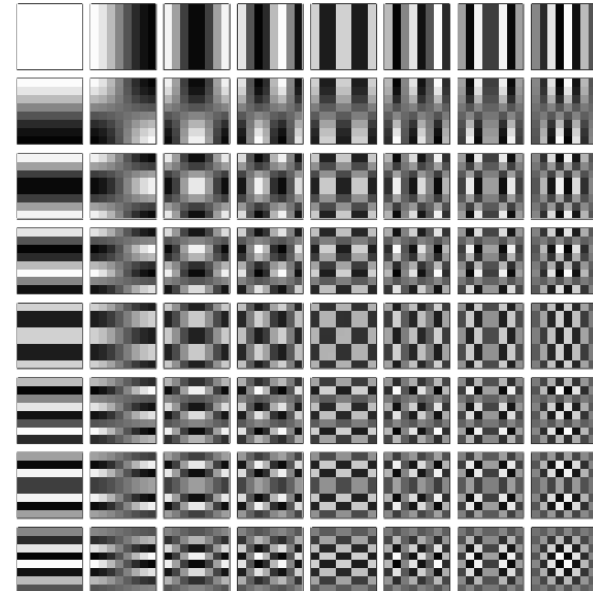
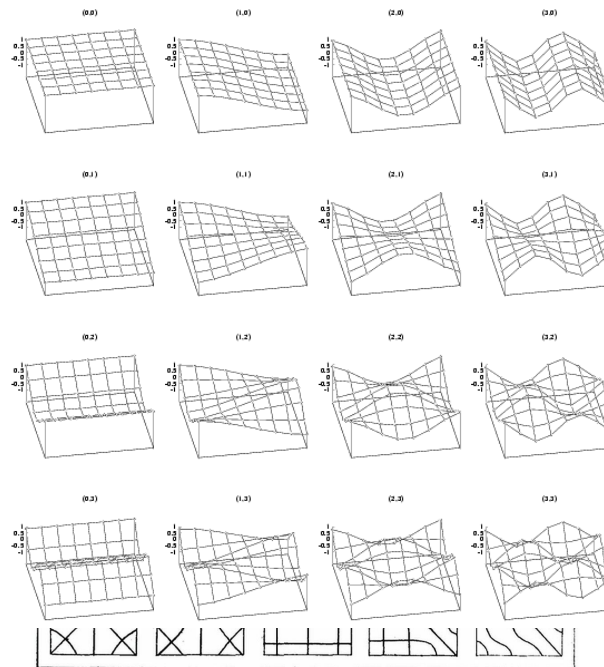
- Other cases where spectrum is known analytically:

- Circle (Bessel functions)
- Cylinder, flat torus (basically rectangle with special bndr. cond.)
- Sphere (Spherical Harmonics)

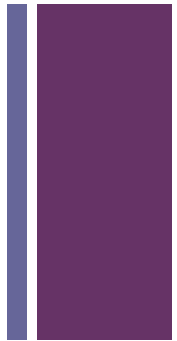


Discrete Cosine Transform

- Similar to Fourier Transform, but only cosines at different frequencies to combine a signal.
- Compression
 - MP3
 - JPEG



+ DCT at Work



8x8

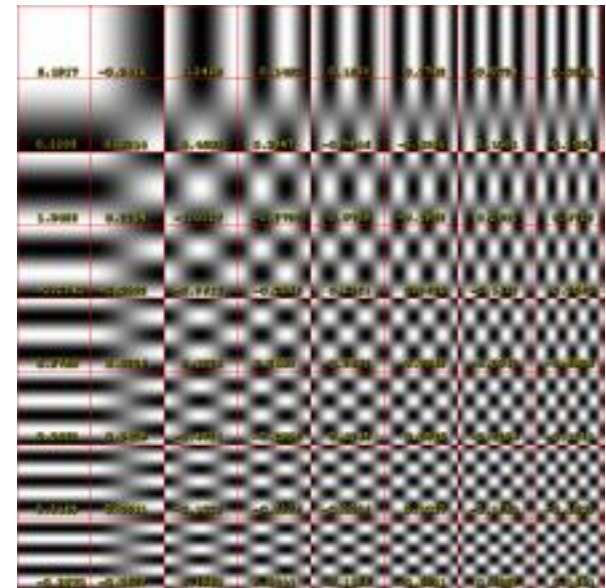
A



6.1917	-0.3411	1.2418	0.1492	0.1583	0.2742	-0.0724	0.0561
0.2205	0.0214	0.4503	0.3947	-0.7846	-0.4391	0.1001	-0.2554
1.0423	0.2214	-1.0017	-0.2720	0.0789	-0.1952	0.2801	0.4713
-0.2340	-0.0392	-0.2617	-0.2866	0.6351	0.3501	-0.1433	0.3550
0.2750	0.0226	0.1229	0.2183	-0.2583	-0.0742	-0.2042	-0.5906
0.0653	0.0428	-0.4721	-0.2905	0.4745	0.2875	-0.0284	-0.1311
0.3169	0.0541	-0.1033	-0.0225	-0.0056	0.1017	-0.1650	-0.1500
-0.2970	-0.0627	0.1960	0.0644	-0.1136	-0.1031	0.1887	0.1444

+

6.192 ✕



Source: Wikipedia

+ Heat Kernel

- Heat Equation:

- Temperature distribution after time t:

$$\frac{\partial}{\partial t} f + \Delta f = 0$$

- Solution (with initial condition

$$f_0): \quad f(\vec{x}, t) = \int_M K(\vec{x}, \vec{y}, t) f_0(\vec{y}) d\vec{y}$$

- Initial condition unit heat source at point p:

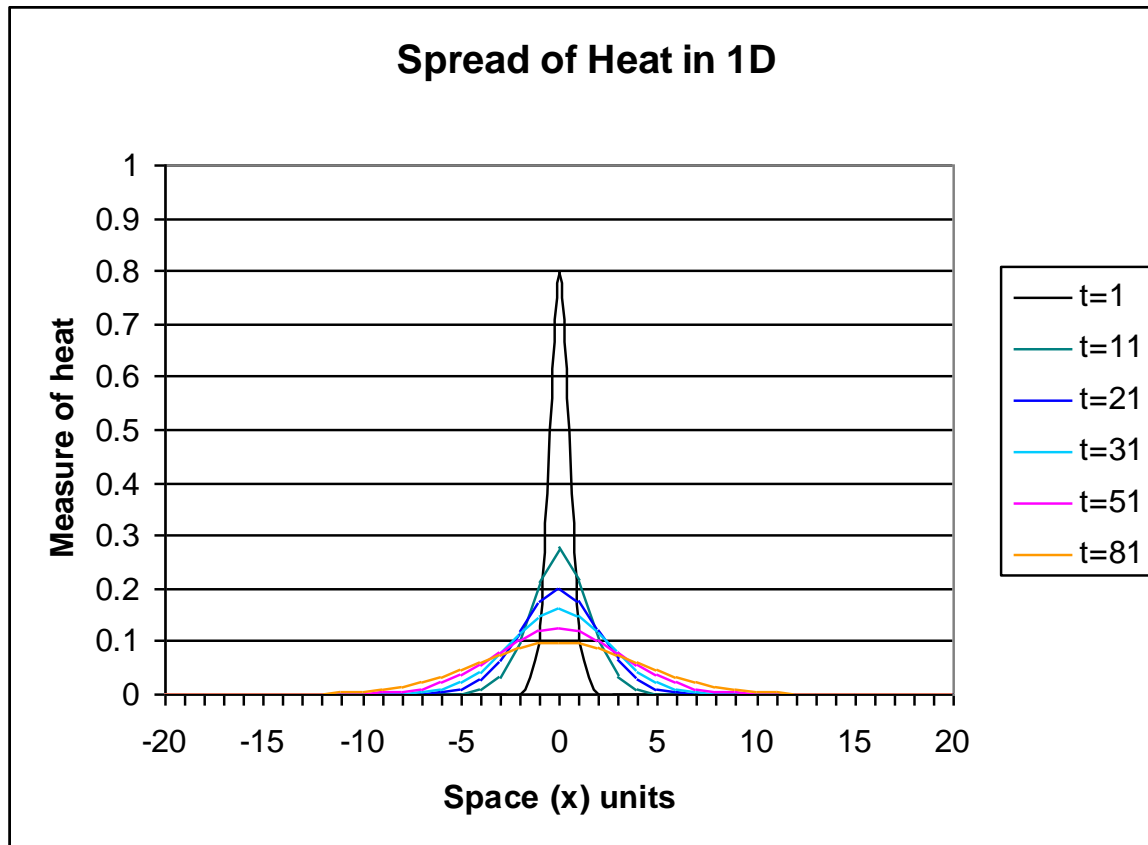
$$f(\vec{x}, t) = K(\vec{x}, \vec{p}, t)$$

- Once we know Eigenfunctions and –values:

$$K(\vec{x}, \vec{y}, t) = \sum_{n=0}^{\infty} \exp^{-\lambda_n t} f_n(\vec{x}) f_n(\vec{y})$$



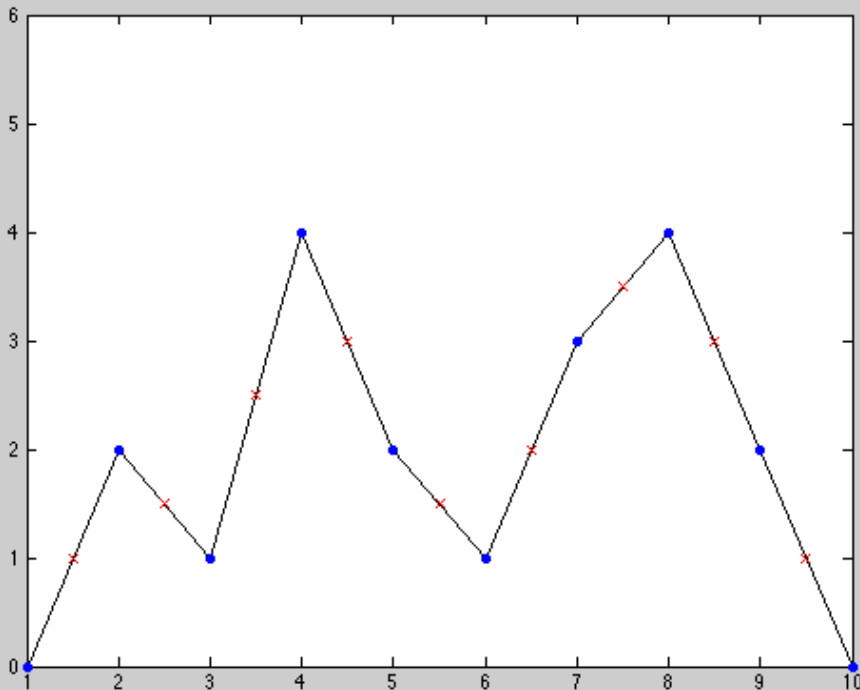
1D Heat Kernel



+ Discrete 1D Laplacian

The Laplace operator is simply the second derivative:

$$\Delta f \coloneqq f'' \quad \text{in } \circ$$



$$\Delta y_i = \left(\frac{y_{i+1} - y_i}{h} - \frac{y_i - y_{i-1}}{h} \right) \frac{1}{h}$$

$$= \frac{y_{i-1} - 2y_i + y_{i+1}}{h^2} = \sum_{j \in \mathbb{N}(i)} \frac{y_i - y_j}{h^2}$$

$$h=1 \Rightarrow \begin{pmatrix} 1 & -2 & 1 \end{pmatrix} \begin{pmatrix} y_{i-1} \\ y_i \\ y_{i+1} \end{pmatrix} \quad (1D \text{ filter})$$

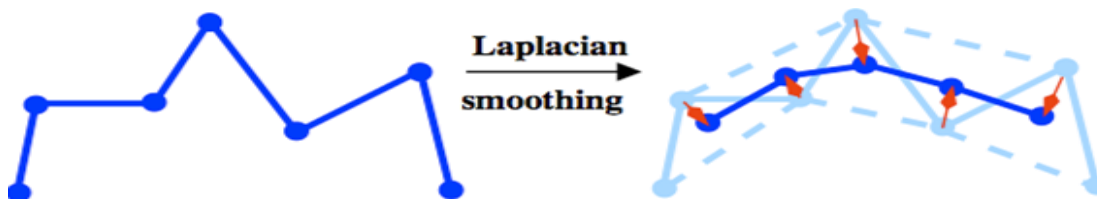
$$\text{Full : } \begin{pmatrix} \dots & -2 & 1 & 0 & 0 \\ 1 & -2 & 1 & 0 & \\ 0 & 1 & -2 & 1 & \\ 0 & 0 & 1 & -2 & \dots \end{pmatrix} y = Ly$$

$$W = \text{adj}(G) \quad \text{and} \quad V_{ii} = \# \text{Neighbors}$$

$$L = W - V$$

+ 1D-Laplacian-Smoothing

Move vertex half way towards center of neighbors:



$$\hat{\mathbf{v}}_i = \frac{1}{2} \left[\frac{1}{2} (\mathbf{v}_{i-1} + \mathbf{v}_i) \right] + \frac{1}{2} \left[\frac{1}{2} (\mathbf{v}_i + \mathbf{v}_{i+1}) \right] = \frac{1}{4} \mathbf{v}_{i-1} + \frac{1}{2} \mathbf{v}_i + \frac{1}{4} \mathbf{v}_{i+1}$$

$$\hat{v}_i = v_i + \frac{1}{4} \Delta v_i \quad \text{where} \quad \Delta v_i = \begin{pmatrix} \Delta x_i \\ \Delta y_i \end{pmatrix}$$

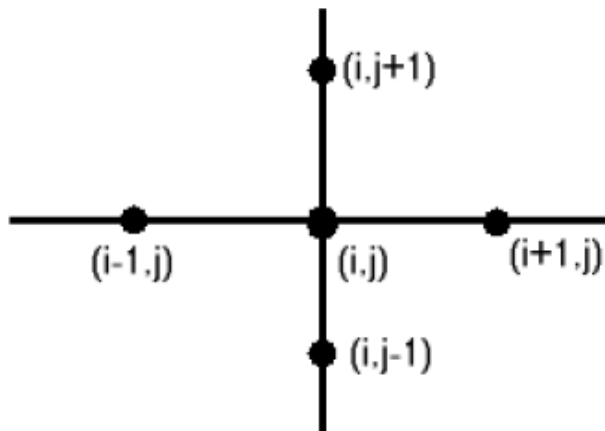
...but don't wash it too long, it will shrink!

+ 2D Grid Laplacian

In 2D Laplace is sum of second partial derivatives:

$$\Delta f(x,y) := f_{xx} + f_{yy} \quad \text{in } \circ^2$$

Discrete 2D-Filter is a 5 Stencil:

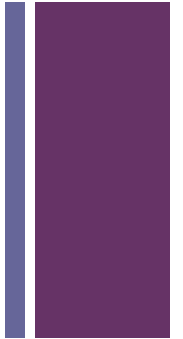


$$\begin{pmatrix} 0 & 1 & 0 \\ 1 & -4 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$

again $L = W - V$



Graph Laplacian



Extension to a Graph is simply (guess):

$$L = W - V \quad \text{or} \quad \Delta y_i = \sum_{j \in \mathcal{N}(i)} (y_j - y_i)$$

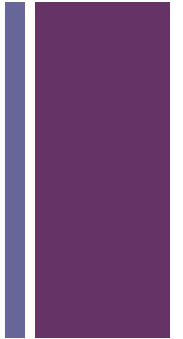
With (symmetric) edge weights (guess again):

$$\Delta y_i = \sum_{j \in \mathcal{N}(i)} w_{ij} (y_j - y_i) \quad \text{or} \quad L := W - V \quad \text{with} \quad W := (w_{ij}) \quad \text{and} \quad V := \text{diag} \left(\sum_{j \in \mathcal{N}(i)} w_{ij} \right)$$

And finally with node weights:

$$\Delta y_i = \frac{1}{d_i} \sum_{j \in \mathcal{N}(i)} w_{ij} (y_j - y_i) \quad \text{or} \quad L := D^{-1} (W - V) \quad \text{with} \quad D := \text{diag}(d_i)$$

+ Spectral Graph Theory

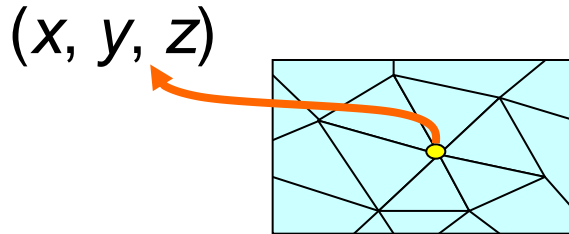


- Eigenvalues
 - Closely related to variety of global graph invariants
 - Global shape descriptors
- Eigenvectors
 - Useful extremal properties, e.g., heuristic for NP-hard problems, normalized cuts and sequencing
 - Spectral embeddings capture global information, e.g., clustering, manifold learning
- However, we'll not look at graphs here, but meshes ...

+ Meshes



Meshes can be seen as graphs:



... with a geometry (here triangle mesh embedded into \mathbb{R}^3).

Therefore one possibility is to use same Laplace discretizations
(with appropriate weights)

How to choose weights “correctly” ?

+ Operator Discretizations

- Graph Laplace:
 - Vertex weights = 1 (masses) and edge weights = 1 (stiffness)
 - Other options (vertex degree, other symmetric weights)
 - Usually not geometry aware
- Mesh Laplace:
 - Different weights (based on geometry)
 - Possibly mesh dependent
 - What if we go to quad-meshes
 - Or tetrahedral meshes (3D)...
- FEM Laplace:
 - Different approach based on Surface or Volume elements
 - Generalizable and higher order approximations

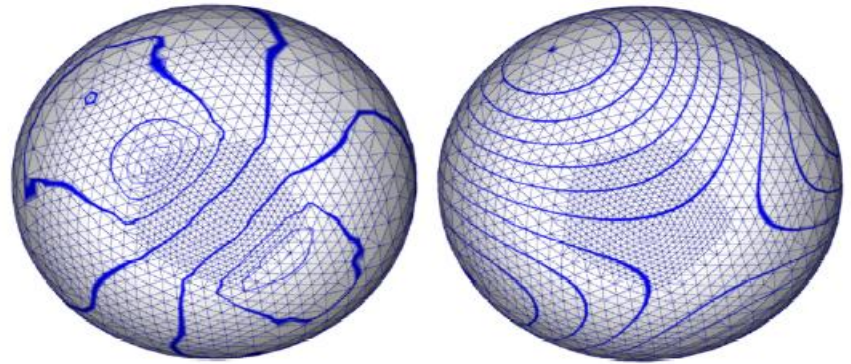


Figure courtesy of Bruno Levy



Mesh Laplacians (Examples)



- Pinkall and Polthier (93): $w_{ij} = \frac{\cot(\alpha_{ij}) + \cot(\beta_{ij})}{2}$ $d_i = 1$

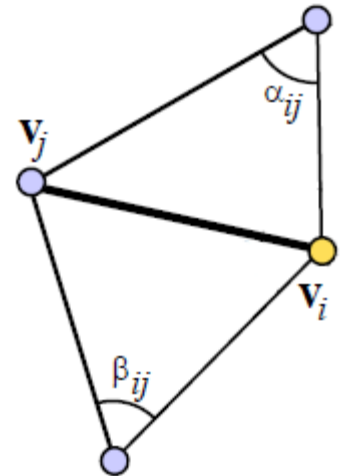
where α and β denote the two angles opposite edge (i,j) . Still depends on mesh sampling (missing mass weights).

- Desbrun (99): $d_i = \frac{area_i}{3}$

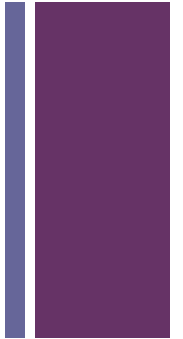
where the area is summed for all triangles at vertex i . Will turn out to be a simplified linear finite element operator.

- Meyer (02): $d_i = \frac{voronoi_i}{3}$

use the the area obtained by joining the circumcenters of the triangles around vertex i (i.e., the Voronoi region).



+ Finite Element Method

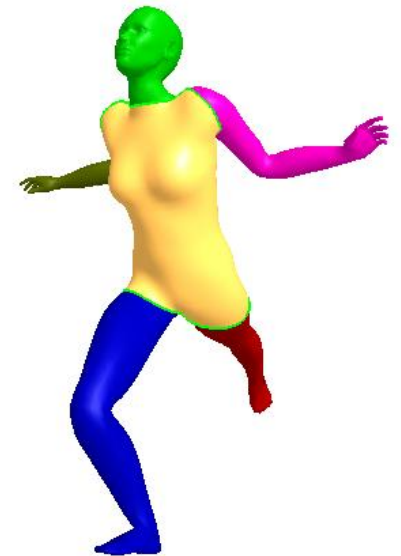
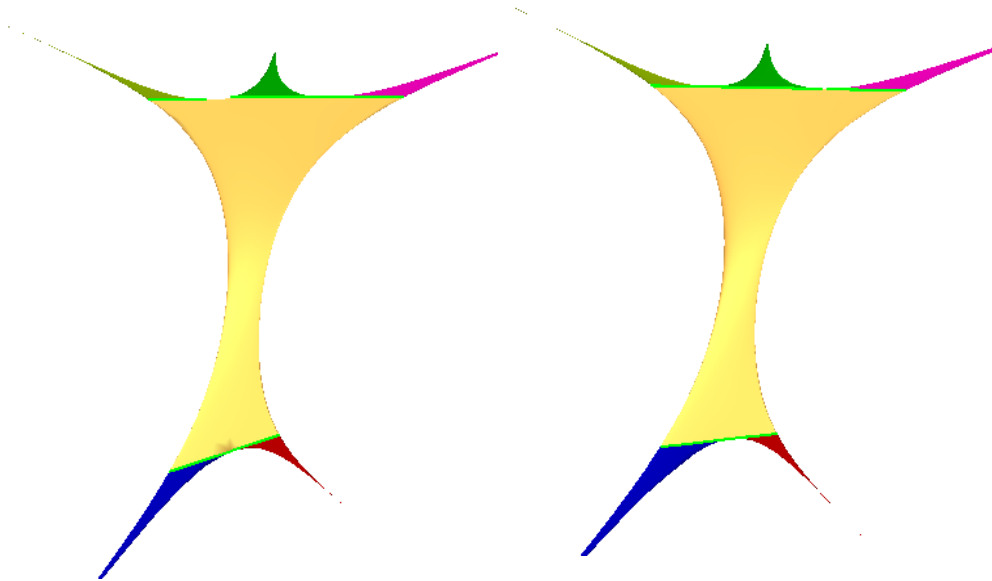


- Instead of graph/wireframe (vertices, edges), we look at elements that assemble our geometry without gaps:
 - triangles
 - tetrahedra
 - voxels....
- We define basis functions over this discretized geometry (linear, quadratic, cubic ...)
- We get a powerful framework to solve differential equations (not just Laplace).
- Details later...

+ Spectral Embedding

- Maps each vertex to the value of all (or a few) Eigenfunctions at that location:

$$\Phi(\vec{x}) = (f_1(\vec{x}), f_2(\vec{x}), f_3(\vec{x}), \dots)$$



+ Spectral Embedding

- Dirichlet Energy measures smoothness
- For Eigenfunctions it is:

$$\begin{aligned} E[f_i] &= \int_M |\nabla f_i|^2 d\sigma = \int_M (\nabla f_i)^2 d\sigma \\ &= - \int_M f_i \Delta f_i d\sigma = \lambda_i \int_M f_i f_i d\sigma = \lambda_i \end{aligned}$$

- Optimal Embedding:
 - The first (non constant) Eigenfunction (also called Fiedler vector) yields the optimal (smoothest) embedding of the shape onto a line (orthogonal to constant function and complying with boundary condition)
 - Higher functions yield smoothest embedding orthogonal to previous ones

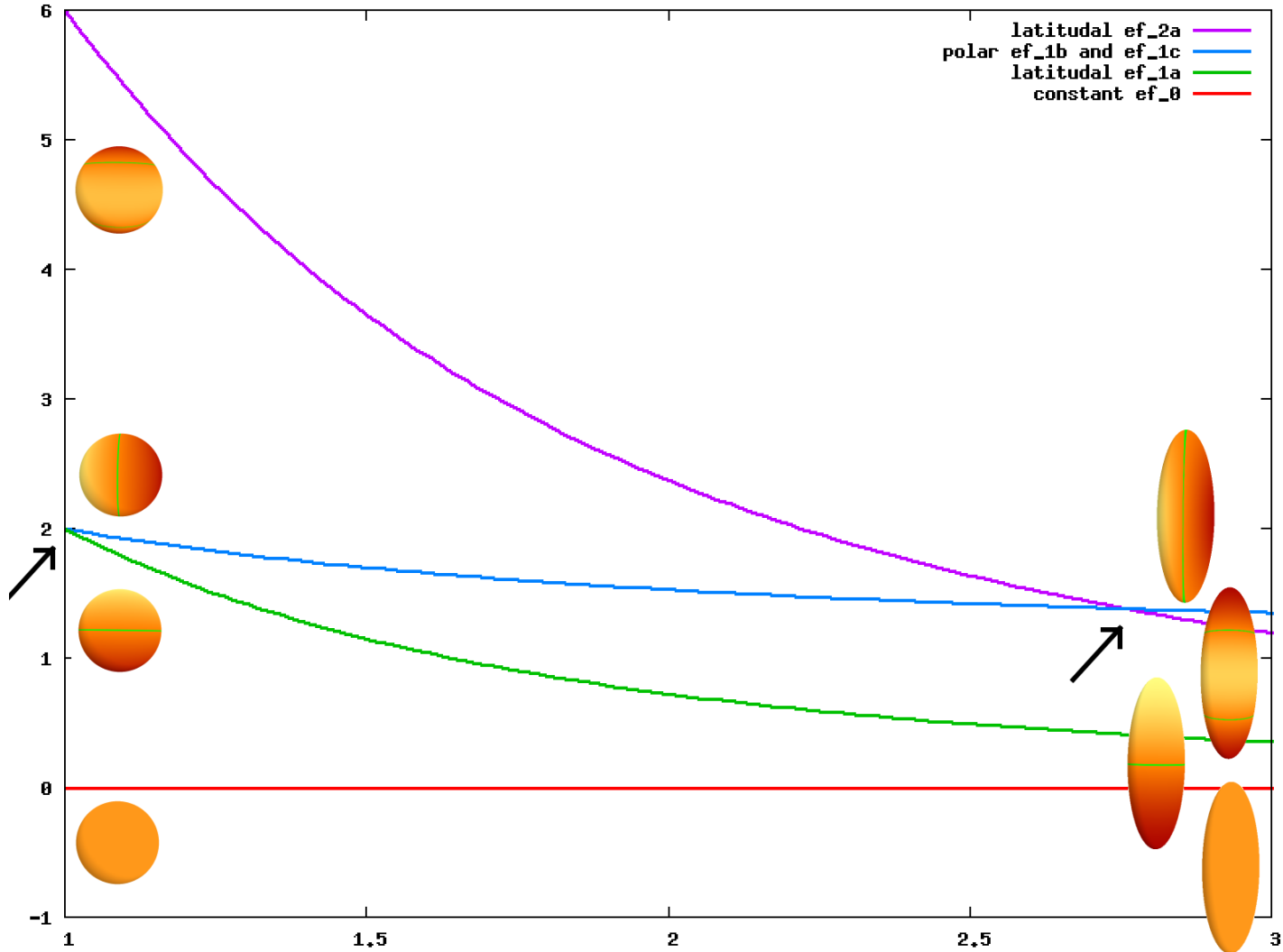


Difficulties

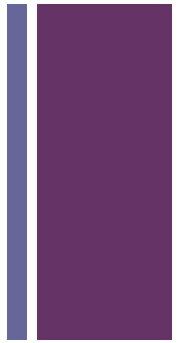


- Numerical inaccuracies
- Sign Flips
 - Scalar multiples are contained in same space
- Higher dimensional Eigenspaces
 - Arbitrary basis (luckily they are rare, but being close to one is already problematic)
 - Due to numerical errors Eigenvalues will be slightly different and we cannot really detect these situations
- Switching of Eigenfunctions
 - Occurs, because of numerical instabilities (two close Eigenvalues switch)
 - Or due to geometric (non-isometric) deformations

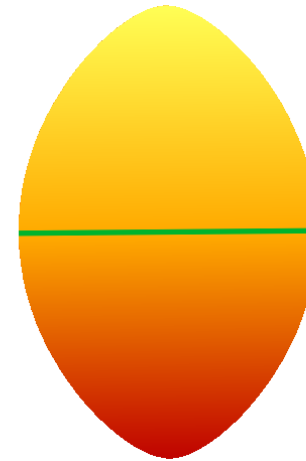
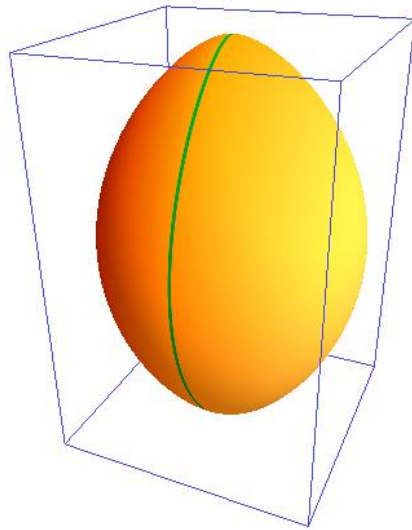
+ Switching of Eigenfunctions



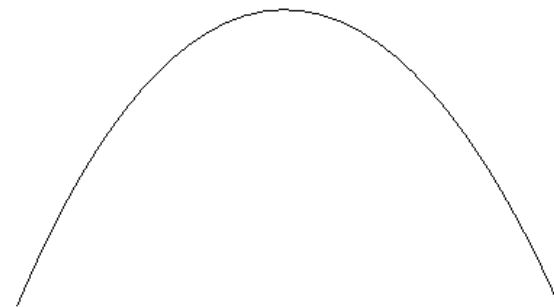
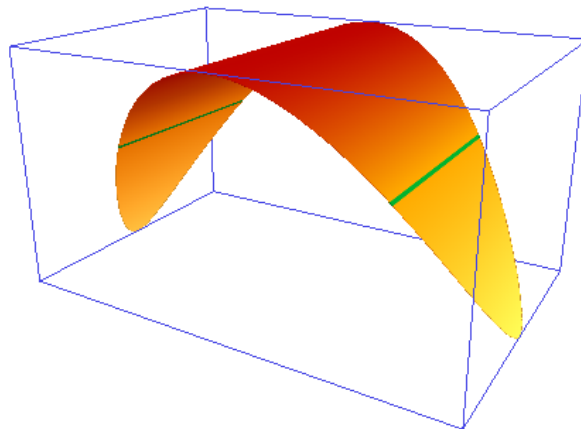
+ Switching of Eigenfunctions



■ $Z=2.74$

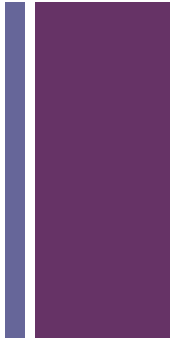


■ $Z=2.76$





Applications: smoothing revisited



- Represent any function as linear combination of Eigenfunctions (similar to DCT):

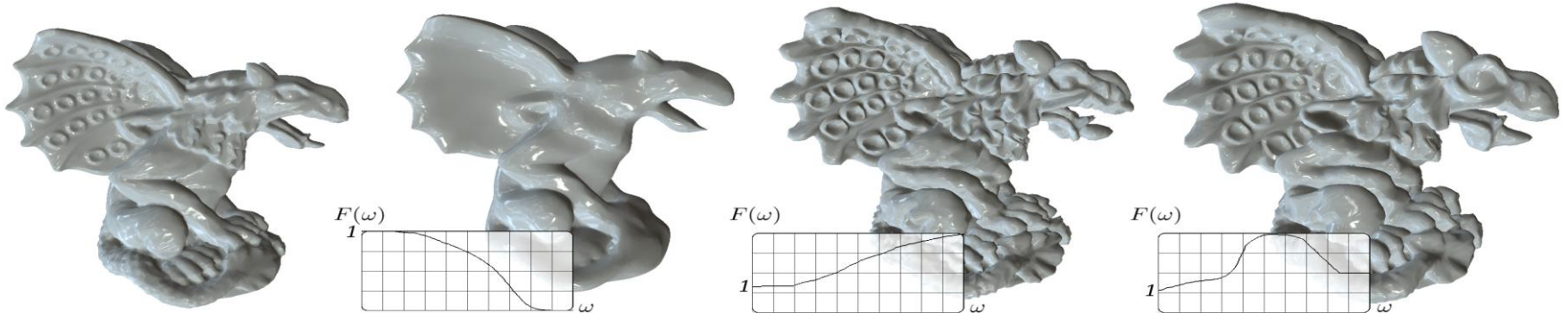
$$F(\vec{x}) = \sum_{i=1}^N c_i f_i$$

- Smooth iteratively (modify coefficients):

$$S^m F = \left(I + \frac{1}{4}L\right)^m F = \sum_{i=1}^n \left(I + \frac{1}{4}L\right)^m c_i f_i = \sum_{i=1}^n \underbrace{\left(1 - \frac{1}{4}\lambda_i\right)^m}_{\text{filter}} c_i f_i$$

+ Geometry Filtering

Take $\mathbf{f} = (x, y, z)$



Courtesy of Bruno Levy

+ Color Filtering on Mesh

Take $\mathbf{f} = (r, g, b)$



Courtesy of Bruno Levy