

## SOLUTIONS TO TEST 2

### SEPARATION OF VARIABLES (estimated time $\sim 20$ min):

1) Consider the two-dimensional Schrödinger equation

$$-\left(\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2}\right) = E\psi, \quad x \in [0, a], \quad y \in [0, b], \quad (1)$$

where  $\psi = \psi(x, y)$  is a wavefunction,  $E$  is the energy, and  $a$  and  $b$  are the dimensions of the rectangle spanned by  $x$  and  $y$ . The wavefunction must satisfy the boundary conditions

$$\frac{\partial \psi(0, y)}{\partial x} = \frac{\partial \psi(a, y)}{\partial x} = 0, \quad \psi(x, 0) = \psi(x, b) = 0, \quad (2)$$

which imposes constraints on the separation constant and the energy  $E$ , for a non-trivial wavefunction to exist, i.e., a wavefunction which is not identically zero. Write  $\psi(x, y) = F(x)G(y)$  and use separation of variables to determine the smallest energy  $E_{\min}$ , for which a non-trivial wavefunction can exist.

### SOLUTION:

1.1) We insert the assumption  $\psi(x, y) = F(x)G(y)$  into Eq. (1) and obtain the equation

$$-F_{xx}G - FG_{yy} = EFG \quad \Rightarrow \quad -\frac{F_{xx}}{F} = \frac{G_{yy}}{G} + E,$$

where subscript denotes differentiation. In the last equation the left hand side depends only on  $x$ , while the right hand side depends only on  $y$ . This can only be satisfied if both sides are equal to the same constant  $\lambda$ . Because we have homogeneous boundary conditions we can then separate into the following two equations

$$\begin{aligned} F_{xx} &= -\lambda F, & F_x(0) &= F_x(a) = 0, \\ G_{yy} &= (\lambda - E)G, & G(0) &= G(b) = 0. \end{aligned}$$

Let us look at the equation for  $F(x)$  first: For  $\lambda=0$  we obtain the solution

$$F_{xx} = 0 \quad \Rightarrow \quad F(x) = c_0 + d_0x \quad \Rightarrow \quad F_x(0) = F_x(a) = d_0 = 0 \quad \Rightarrow \quad F(x) = c_0.$$

Thus  $\lambda=0$  is an eigenvalue with eigenfunction  $c_0$ . For  $\lambda = -k^2 < 0$  we obtain the solution

$$F(x) = c_1 e^{kx} + d_1 e^{-kx}.$$

The left boundary condition is  $F_x(0) = k(c_1 - d_1) = 0$ . Thus  $F(x) = 2c_1 \cosh(kx)$  and the right boundary condition  $F_x(a) = 2c_1 k \sinh(ka) = 0$  cannot be fulfilled unless  $a=0$  (unphysical box) or  $c_1 = d_1 = 0$  (trivial zero-solution). Therefore  $\lambda < 0$  is also not an eigenvalue.

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For  $\lambda = k^2 > 0$  we obtain the solution

$$F(x) = c_2 \cos(kx) + d_2 \sin(kx).$$

The left boundary condition gives

$$F_x(0) = kd_2 = 0 \quad \Rightarrow \quad d_2 = 0 \quad \Rightarrow \quad F(x) = c_2 \cos(kx).$$

To obtain a nontrivial solution we must therefore require that  $c_2 \neq 0$  and again  $a > 0$  to have a physical box size. The right boundary condition then gives us the eigenvalues

$$F_x(a) = -kc_2 \sin(ka) = 0 \quad \Rightarrow \quad \lambda = \lambda_n = k_n^2 = \left(\frac{n\pi}{a}\right)^2, \quad n = 1, 2, 3 \dots$$

Now we remember that  $\lambda=0$  was an eigenvalue with the eigenfunction being a constant. This case is contained in the above expression, and thus we may combine them and write the final solution

$$F = F_n(x) = A_n \cos(\sqrt{\lambda_n} x), \quad \lambda_n = \left(\frac{n\pi}{a}\right)^2, \quad n = 0, 1, 2, 3 \dots$$

We can now write the equation for  $G(y)$  in the form

$$G_{yy} = -\gamma G, \quad G(0) = G(b) = 0 \quad \gamma = E - \lambda_n.$$

For  $\gamma=0$  the solution is  $G(y)=c_3 + d_3y$  and the boundary conditions  $G(0)=c_3=0$  and  $G(b)=d_3b=0$  can only be fulfilled by the trivial zero-solution  $c_3=d_3=0$  or an unphysical box size  $b=0$ . Thus  $\gamma=0$  is not an eigenvalue.

For  $\gamma=-\kappa^2 < 0$  the solution is  $G(y)=c_4 \exp(\kappa y) + d_4 \exp(-\kappa y)$ . The boundary condition  $G(0)=c_4 + d_4=0$  gives  $G(y)=2c_4 \sinh(\kappa y)$  and therefore  $G(b) = 2c_4 \sinh(\kappa b) = 0$  can only be fulfilled by the trivial zero-solution  $c_4=d_4=0$  or an unphysical box size  $b=0$ . Thus  $\gamma < 0$  is also not an eigenvalue.

Finally, for  $\gamma = \kappa^2 > 0$  we obtain the solution  $G(y) = c_5 \cos(\kappa y) + d_5 \sin(\kappa y)$ . The left boundary condition gives  $G(0)=c_5=0$  and thus  $G(y)=d_5 \sin(\kappa y)$ . To obtain a nontrivial solution we must therefore require that  $d_5 \neq 0$  and  $b > 0$  to have a physical box size. The right boundary condition then gives us the eigenvalues

$$G(b) = d_5 \sin(\kappa b) = 0 \quad \Rightarrow \quad \gamma = \gamma_m = \kappa_m^2 = \left(\frac{m\pi}{b}\right)^2, \quad m = 1, 2, 3 \dots$$

We can now determine the minimum energy from the relation  $E=\gamma_m + \lambda_n$  as follows

$$E = \gamma_m + \lambda_n = \left(\frac{m\pi}{b}\right)^2 + \left(\frac{n\pi}{a}\right)^2 \geq \left(\frac{\pi}{b}\right)^2 = E_{\min},$$

because  $n \geq 0$  and  $m \geq 1$ . Note how it was not necessary to write up the full solution to find the minimum energy.

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### GREEN'S FUNCTION (estimated time $\sim 15$ min.):

2) Consider the following equation

$$\frac{d^2 u}{dx^2} = f(x), \quad x \in [0, 1], \quad u(0) = u(1) = 0, \quad (3)$$

for the function  $u=u(x)$ . Find Green's function  $G(x, x')$  for the problem (3).

#### SOLUTION:

2.1) First we note that the boundary conditions are homogeneous and thus we may apply Green's function approach directly with the same homogeneous boundary conditions. We define Green's function  $G(x, x')$  to be a solution to the problem

$$\frac{d^2 G}{dx^2} = \delta(x - x'), \quad G(0, x') = G(1, x') = 0, \quad x \in [0, 1], \quad x' \in ]0, 1[ \quad (4)$$

For  $0 \leq x < x'$  Green's function satisfies the homogeneous equation  $d^2 G/dx^2 = 0$  and we obtain the solution

$$G = G_1(x, x') = Ax + B \quad \Rightarrow \quad G_1(0, x') = B = 0 \quad \Rightarrow \quad G_1(x, x') = Ax.$$

For  $x' < x \leq 1$  Green's function satisfies also the homogeneous equation  $d^2 G/dx^2 = 0$  and we obtain the solution

$$G = G_2(x, x') = C(x - 1) + D \quad \Rightarrow \quad G_2(1, x') = D = 0 \quad \Rightarrow \quad G_2(x, x') = C(x - 1).$$

Integrating Eq. (4) from  $x' - \epsilon$  to  $x' + \epsilon$  and letting  $\epsilon \rightarrow 0$  we obtain the condition

$$\left. \frac{dG_2}{dx} \right|_{x=x'} - \left. \frac{dG_1}{dx} \right|_{x=x'} = C - A = 1 \quad \Rightarrow \quad C = A + 1.$$

From Eq. (4) we see that  $d^2 G/dx^2$  has a  $\delta$ -function in  $x'$  and that  $dG/dx$  has a discontinuity in  $x'$ . This means that Green's function itself is continuous in  $x'$ , which gives the requirement

$$G_2(x', x') = G_1(x', x') \quad \Rightarrow \quad Ax' = C(x' - 1) = (A + 1)(x' - 1) \quad \Rightarrow \quad A = x' - 1.$$

The expression for Green's function is therefore

$$G(x, x') = \begin{cases} (x' - 1)x & \text{for } 0 \leq x < x' \\ (x - 1)x' & \text{for } x' < x \leq 1 \end{cases}$$

As a check we note that Green's function  $G(x, x')$  is symmetric in its arguments, as it should be. We know that this should be the case because the operator  $\hat{L} = d^2/dx^2$  appearing in Eq. (3) is a Sturm-Liouville operator and Green's function is real.